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# 双排抗滑桩加固滑坡的前桩后侧推力算法

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**摘要:**为了在求解双排全长桩和后排沉埋-前排全长桩中前桩后侧坡体推力的同时也可得到其分布模式,以便更合理确定 前排桩受力特征,针对双排桩加固的基岩-覆盖层式滑坡,采用斜条分法将排间滑体分割成若干斜向平行滑面的土条;对所 形成的在不同深度位置的4类典型土条,根据静力平衡条件,推导了其作用于前桩受荷段后侧的推力计算公式,进而可确 定出前排桩受荷段后侧坡体推力分布模式及其合力。实例分析表明,本文方法所确定的前桩后侧坡体推力呈不规则抛物 线形分布,其峰值点接近于滑面;后排桩沉埋模式时前桩后侧坡体推力比后排桩全长模式时的值小8.6%~10.6%;本文方法 比传递系数法的计算值偏大10.6%~13.4%,且相对更接近于FLAC<sup>3D</sup>模拟结果。 关键词:双排抗滑桩;滑坡推力;沉埋模式;斜条分法;排间滑体

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## A calculation method for thrust on the fore piles of double-row stabilizing piles used to reinforce landslides

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**Abstract:** In order to better understand the thrust and its distribution pattern on the fore piles of double-row fulllength piles and embedded rear-row piles combined with full-length fore-row ones used to reinforce landslides, an inclined slice method is provided in this article. The slide mass between the two pile-rows is divided into many oblique slices parallel to the local slip surface. These slices are involved in four basic types at various depths. Calculation formula of the thrust of each inclined slice on the loading segment of the fore pile can be derived based on the static equilibrium principle, and the distribution pattern of the thrust and its resultant force can be determined. The analysis results of the examples indicate that the thrust on the fore pile determined by the proposed method shows approximately parabolic distribution and its peak point is close to the slip surface. The calculated thrust on the fore piles under the embedded rear-row piles are 8.6% to 10.6% smaller than that under the full-length rear-row piles. The proposed thrusts are about 10.6% to 13.4% higher than those by the classical transfer coefficient method, and are relatively close to the numerical results via FLAC<sup>3D</sup>.

**Keywords**: double-row stabilizing piles; landslide thrust; embedded rear-row piles layout; inclined slice method; slide mass between two pile-rows

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双排全长桩或后排沉埋-前排全长桩越来越多应 用于大型基覆式滑坡的治理工程<sup>[1-6]</sup>,如图1所示, *d*<sub>0</sub>=0与*d*<sub>0</sub>>0分别表示前排全长与后排沉埋桩。对于 此类问题,两排桩上推力大小及分布模式是工程设计 中的关键因素之一。



Fig. 1 Sketch map of the cross section of a landslide reinforced with double-row piles

后排桩上推力求解类似于单排桩,以往已有较多 计算方法<sup>[7-13]</sup>。中铁二院曾对单排桩上推力进行现场 模型试验<sup>[14-16]</sup>,得到了其分布模式。也有学者通过室 内模型试验对沉埋桩及后排沉埋-前排全长桩后侧推 力进行研究<sup>[17-19]</sup>,确定了试验模型桩后推力作用特 征。郑颖人等<sup>[20]</sup>、赵尚毅等<sup>[21]</sup>、宋雅坤等<sup>[22]</sup>则采用有 限元方法得到了单排全长抗滑桩或沉埋桩的受力特征。

在前排桩后侧推力的解析计算方面,刘鸿<sup>[23]</sup>、李 明康<sup>[24]</sup>采用弹性理论对双排抗滑桩排间岩土体的传 力机制进行了分析,得出了前排桩后侧推力的表达 式,但对排间岩土体,其采用弹性分析模型偏于理想 化;肖世国等<sup>[25]</sup>基于后排桩推力向前传递的地基系数 "k"法以及排间岩土体推力传递的传递系数法,提出了 一种计算前排桩后侧推力的方法,但不能给出前排桩 上推力分布模式,且采用滑体侧向抗力系数沿深度不 变的模式也尚需优化;杨磊<sup>[26]</sup>将弹性理论与条分法相 结合,以此可近似得到前排桩上推力,但该法关于条 块间正应力分布模式呈抛物线形或直线型的假设仍 存在不合理性。总之,上述这些求解前排桩上推力的 方法有下面两个明显缺陷:(1)模型过于理想化;(2)无 法兼顾推力大小与分布模式。

针对既有方法的不足,为了在求解前排桩后侧坡 体推力的同时也可得到其分布模式,以便更合理确定 前排桩受力特征,本文采取如下基本思路:首先,考虑 滑体侧向抗力系数随深度变化的实际情况,采用地基 系数"*m*"法求得后排桩受荷段前侧滑体抗力;然后,考 虑双排桩排间滑面特征,采用平行于滑面的斜条分法 对排间土体进行分析,将排间土体分割为若干斜条, 根据各条块应满足的静力平衡条件,同时确定作用于 前排桩受荷段后侧的推力大小及其分布模式。

## 1 计算方法

对于双排桩加固滑坡的排间滑体,其基本分析模型如图 2 所示,将其后边界 OC 与前边界 DE 之间的滑体以平行排间滑面方向分割成 n+1 条,边界 OC 与 DE 自上而下的分割点分别依次标为点 A<sub>0</sub>, A<sub>1</sub>, …, A<sub>n-1</sub> 和 点 D<sub>1</sub>, D<sub>2</sub>, …, D<sub>n-1</sub>。图 2 中,  $d_0$  为后排桩沉埋深度,当  $d_0=0$  时,表示全长桩; d 为后排桩受荷段长度; s 为前后排间距;  $s_0$  为排间滑面折点 F 距后排桩距离;  $q_1$ 、 $q_1+q_2$ 分别为边界 OC 顶、底端的水平压力;  $a_1$ 、 $a_2$ 分别为 排间滑面段 CF、FE 的水平倾角;  $e_1$ 、i、n 为条块编号。



对于排间滑面 CE,这里考虑存在一个折点 F 的情况。当 α<sub>1</sub>=α<sub>2</sub> 时,退化为直线型滑面;对于排间滑面多 折点的情况,求解方法类似于一个折点情况,不再赘 述。为简化分析,作如下假定:

(1)后排桩受荷段前侧土体的抗力在后排桩全长 与沉埋时,分别为三角形、两段折线形分布。

(2)各斜条块假设为刚体,且条块之间相互作用 力、前排桩作用于各条块的力、滑床对最下面一个条 块的作用力均位于相应作用面的中点,由此可得各典 型条块受力分析模型如图 3 所示。

(3)前排桩作用于第 *i* 条块的切向力 *T*<sub>pi</sub>和法向力 *N*<sub>pi</sub>之间满足 *T*<sub>pi</sub>=*k*<sub>p</sub>*N*<sub>pi</sub>, 其中 *k*<sub>p</sub>为桩土界面综合参数。

(4)两排桩之间的滑面处于极限状态且满足 Coulomb 强度准则。

记排间滑体重度、内摩擦角、黏聚力分别为γ、

φ、c, 对图 3 所示的不同深度处典型条块分别进行受力分析。









图 3 各典型斜条块受力分析模型 Fig. 3 Mechanical analysis model of each typical oblique slice

1.1 顶部条块 e<sub>1</sub>

对于三角形条块 *e*<sub>1</sub>,只包含其下表面上的切向力 *T*<sub>0</sub>和法向力 *N*<sub>0</sub> 共 2 个未知力,需满足水平和竖向静 力平衡方程,即:

$$\begin{cases} \frac{q_1 h_0^2}{2d_0} + N_0 \sin \alpha_2 - T_0 \cos \alpha_2 = 0\\ G_{e_1} - N_0 \cos \alpha_2 - T_0 \sin \alpha_2 = 0 \end{cases}$$
(1)

$$G_0$$
 — 条块  $e_1$  的重力,  $G_0 = \gamma s h_0/2$ ;  
 $l_0$  — 条块  $e_1$  底面长度。  
根据 Coulomb 强度准则, 还应满足:

 $T_0 \leq N_0 \tan \varphi + cl_0 \tag{2}$ 

$$\begin{cases} N_0 = \frac{\gamma s h_0}{2} \cos \alpha_2 - \frac{q_1 h_0^2}{2d_0} \sin \alpha_2 \\ T_0 = \frac{\gamma s h_0}{2} \sin \alpha_2 + \frac{q_1 h_0^2}{2d_0} \cos \alpha_2 \end{cases}$$
(3)

## 1.2 位于 d<sub>0</sub> 范围的平行四边形条块

此范围内各条块分析方法相同,以条块1为例具体阐述。以条块1下表面上的切向力 N<sub>1</sub>和法向力 T<sub>1</sub>作用点位置为矩心,由静力平衡条件及假定(3)可得:

$$\begin{cases} G_{1} - T_{1} \sin \alpha_{2} - N_{1} \cos \alpha_{2} + T_{0} \sin \alpha_{2} + N_{0} \cos \alpha_{2} - T_{p1} = 0\\ Q_{1} + N_{1} \sin \alpha_{2} - T_{1} \cos \alpha_{2} - N_{0} \sin \alpha_{2} + T_{0} \cos \alpha_{2} - N_{p1} = 0\\ M_{1} - N_{0}h_{1} \sin \alpha_{2} + T_{0}h_{1} \cos \alpha_{2} + \frac{1}{2}N_{p1}(s \tan \alpha_{2} - h_{1})\\ - \frac{1}{2}T_{p1}s = 0\\ T_{p1} = k_{p}N_{p1} \end{cases}$$
(4)

$$\begin{cases} Q_1 = \frac{q_1 (2h_0 + h_1) h_1}{2d_0} \\ M_1 = \int_{x_{A_0}}^{x_{A_1}} \frac{q_1 x}{d_0} \left( h_0 + h_1 + \frac{s \tan \alpha_2}{2} - x \right) dx \\ G_1 = \gamma h_1 s \end{cases}$$
(5)

式中: h1 — 条块1 左侧长度, 条块2~n-1 标记以此 类推;下文中 Ni、Ti 的含义同条块1。 由式(4)可推导得到:

$$\left[ \frac{G_{1}-k_{p}\frac{2M_{1}+2T_{0}h_{1}\cos\alpha_{2}-2N_{0}h_{1}\sin\alpha_{2}}{k_{p}s-s\tan\alpha_{2}+h_{1}}}{\cos\alpha_{2}} + \frac{Q_{1}-\frac{2M_{1}+2T_{0}h_{1}\cos\alpha_{2}-2N_{0}h_{1}\sin\alpha_{2}}{k_{p}s-s\tan\alpha_{2}+h_{1}}}{\sin\alpha_{2}} + \frac{Q_{1}-\frac{2M_{1}+2T_{0}h_{1}\cos\alpha_{2}-2N_{0}h_{1}\sin\alpha_{2}}{k_{p}s-s\tan\alpha_{2}+h_{1}}}{(k_{p}s-s\tan\alpha_{2}+h_{1})\sin\alpha_{2}} + \frac{Q_{1}-\frac{2M_{1}+2T_{0}h_{1}\cos\alpha_{2}-2N_{0}h_{1}\sin\alpha_{2}}{(k_{p}s-s\tan\alpha_{2}+h_{1})}}{(k_{p}s-s\tan\alpha_{2}+h_{1})\sin\alpha_{2}} + \frac{Q_{1}-\frac{2M_{1}+2T_{0}h_{1}\cos\alpha_{2}-2N_{0}h_{1}\sin\alpha_{2}}{(k_{p}s-s\tan\alpha_{2}+h_{1})\sin\alpha_{2}}}{(k_{p}s-s\tan\alpha_{2}+h_{1})\sin\alpha_{2}} + \frac{Q_{1}-\frac{2M_{1}+2T_{0}h_{1}\cos\alpha_{2}-2N_{0}h_{1}\sin\alpha_{2}}{(k_{p}s-s\tan\alpha_{2}+h_{1})\sin\alpha_{2}}}{(k_{p}s-s\tan\alpha_{2}+h_{1})} \right] \cot\alpha_{2}$$

对位于 d<sub>0</sub> 范围内的其他条块 k,只需要将方程(6) 中的下标 1 换为 k,将 0 换为 k-1 即可,而其中的 Q<sub>k</sub>、 M<sub>k</sub>、G<sub>k</sub>则按方程(7)求解:

$$\begin{cases} Q_k = \int_{x_{Ak-1}}^{x_{Ak}} \frac{q_1 x}{d_0} dx \\ M_k = \int_{x_{Ak-1}}^{x_{Ak}} \frac{q_1 x}{d_0} \left( x_{Ak} + \frac{s \tan \alpha_2}{2} - x \right) dx \\ G_k = \gamma h_k s \end{cases}$$
(7)

1.3 位于 d 范围内的平行四边形条块

对位于 d 范围内的条块 i, 其受力分析模式同条 块 1, 因此只需将方程(6)中下标 1 换为 i, 将 0 换为 i-1

即可。但该范围内 Q<sub>i</sub>、M<sub>i</sub>解法因左边界荷载不同而与 d<sub>0</sub> 范围内情况略有不同,其求解方程为:

$$\begin{cases} Q_{i} = \int_{x_{Ai-1}}^{x_{Ai}} \left[ \frac{q_{2} - q_{1}}{d_{0}} \left( x - d_{0} \right) + q_{1} \right] dx \\ M_{i} = \int_{x_{Ai-1}}^{x_{Ai}} \left[ \frac{q_{2} - q_{1}}{d_{0}} \left( x - d_{0} \right) + q_{1} \right] \left( x_{Ai} + \frac{s \tan \alpha_{2}}{2} - x \right) dx \end{cases}$$
(8)

1.4 底部条块 n

设该条块面积为 $S_n$ ,则其重力为 $G_n = \gamma S_n$ ;以 $T_{pn}$ 和 $N_{pn}$ 作用点位置为矩心,由静力平衡条件及假定(4)可得:

$$\begin{cases} N_{n1}\sin\alpha_{1} - T_{n1}\cos\alpha_{1} + N_{n2}\sin\alpha_{2} - T_{n2}\cos\alpha_{2} - N_{n-1}\sin\alpha_{2} + T_{n-1}\cos\alpha_{2} - N_{pn} = 0 \\ -N_{n1}\cos\alpha_{1} - T_{n1}\sin\alpha_{1} - N_{n2}\cos\alpha_{2} - T_{n2}\sin\alpha_{2} + N_{n-1}\cos\alpha_{2} + T_{n-1}\sin\alpha_{2} + G_{n} - T_{pn} = 0 \\ N_{n1}\left\{\frac{\left(s - \frac{s_{0}}{2}\right)}{\cos\alpha_{1}} - \left[\frac{\left(s - s_{0}\right)\sin\left(\alpha_{1} - \alpha_{2}\right)}{\cos\alpha_{2}\sin\left(90^{\circ} - \alpha_{1}\right)} + \frac{h_{n}}{2}\right]\sin\alpha_{1}\right\} + T_{n1}\left[\frac{\left(s - s_{0}\right)\sin\left(\alpha_{1} - \alpha_{2}\right)}{\cos\alpha_{2}\sin\left(90^{\circ} - \alpha_{1}\right)} + \frac{h_{n}}{2}\right]\cos\alpha_{1} + \\ N_{n2}\left[\frac{\left(s - s_{0}\right)}{2\cos\alpha_{2}} - \frac{h_{n}}{2}\sin\alpha_{2}\right] + T_{n2}\frac{h_{n}}{2}\cos\alpha_{2} - N_{n-1}\left[\frac{s}{2\cos\alpha_{2}} + \frac{h_{n}}{2}\sin\alpha_{2}\right] + T_{n-1}\frac{h_{n}}{2}\cos\alpha_{2} - G_{n}s_{G} = 0 \\ T_{n1} = N_{n1}\tan\varphi + cl_{CF} \\ T_{n2} = N_{n2}\tan\varphi + cl_{FE} \end{cases}$$

$$(9)$$

式中: s<sub>G</sub> — 条块 n 的形心位置距右侧的水平距离,

$$s_{\rm G} = \frac{s_0^2 - 3ss_0 + 3s^2}{3(2s - s_0)};$$
  
N<sub>n1</sub>、T<sub>n1</sub>——条块 n 底面 CF 上的法向力和切向力;

*N<sub>n2</sub>、T<sub>n2</sub>*—条块*n*底面FE上的法向力和切向力;
 *h<sub>n</sub>*—条块*n*右侧长度;
 *l*<sub>CF</sub>、*l<sub>FE</sub>*—底面CF、FE的长度。
 由式(9)可得关于*N<sub>n2</sub>、N<sub>pn</sub>*的二元一次方程组为:

$$\begin{cases} N_{n2} \left( \frac{\eta_{3}}{\eta_{1}} - \eta_{2} \right) + \eta_{4} + \frac{T_{n-1} \sin \alpha_{2} + N_{n-1} \cos \alpha_{2}}{(\cos \alpha_{1} + \tan \varphi \sin \alpha_{1})} + \frac{G_{n} - k_{p} N_{pn}}{(\cos \alpha_{1} + \tan \varphi \sin \alpha_{1})} + \frac{T_{n-1} \frac{h_{n}}{2} \cos \alpha_{2} - N_{n-1} \left[ \frac{s}{2 \cos \alpha_{2}} + \frac{h_{n}}{2} \sin \alpha_{2} \right] - G_{n} s_{G}}{\eta_{1}} = 0$$

$$\begin{cases} N_{n2} - \chi_{2} N_{pn} - \chi_{3} + \frac{T_{n-1} \cos \alpha_{2} - N_{n-1} \sin \alpha_{2}}{(\sin \alpha_{1} - \tan \varphi \cos \alpha_{1}) \chi_{1}} + \frac{T_{n-1} \sin \alpha_{2} + N_{n-1} \cos \alpha_{2}}{(\cos \alpha_{1} + \tan \varphi \sin \alpha_{1}) \chi_{1}} + \frac{G_{n} - k_{p} N_{pn}}{(\cos \alpha_{1} + \tan \varphi \sin \alpha_{1}) \chi_{1}} = 0 \end{cases}$$

$$(10)$$

式中中间变量  $\eta_1$ 、 $\eta_2$ 、 $\eta_3$ 、 $\eta_4$  和  $\chi_1$ 、 $\chi_2$ 、 $\chi_3$ 、 $\chi_4$  表达式分别为:

$$\begin{pmatrix} \eta_1 = \left\{ \frac{\left(s - \frac{s_0}{2}\right)}{\cos \alpha_1} - \left[ \frac{\left(s - s_0\right) \sin(\alpha_1 - \alpha_2)}{\cos \alpha_2 \sin(90^\circ - \alpha_1)} + \frac{h_n}{2} \right] \sin \alpha_1 \right\} + \left[ \frac{\left(s - s_0\right) \sin(\alpha_1 - \alpha_2)}{\cos \alpha_2 \sin(90^\circ - \alpha_1)} + \frac{h_n}{2} \right] \tan \varphi \cos \alpha_1$$

$$\eta_2 = \frac{\cos \alpha_2 + \tan \varphi \sin \alpha_2}{\cos \alpha_1 + \tan \varphi \sin \alpha_1}$$

$$\eta_3 = \left[ \frac{\left(s - s_0\right)}{2 \cos \alpha_2} - \frac{h_n}{2} \sin \alpha_2 \right] + \frac{h_n}{2} \tan \varphi \cos \alpha_2$$

$$\eta_4 = \frac{cl_{CF} \left[ \frac{\left(s - s_0\right) \sin(\alpha_1 - \alpha_2)}{\cos \alpha_2 \sin(90^\circ - \alpha_1)} + \frac{h_n}{2} \right] \cos \alpha_1 + cl_{FE} \frac{h_n}{2} \cos \alpha_2}{\eta_1} - \frac{cl_{CF} \sin \alpha_1 + cl_{FE} \sin \alpha_2}{\cos \alpha_1 + \tan \varphi \sin \alpha_1}$$

$$(11)$$

$$\begin{cases} \chi_{1} = \frac{\sin \alpha_{2} - \tan \varphi \cos \alpha_{2}}{\sin \alpha_{1} - \tan \varphi \cos \alpha_{1}} - \frac{\cos \alpha_{2} + \tan \varphi \sin \alpha_{2}}{\cos \alpha_{1} + \tan \varphi \sin \alpha_{1}} \\ \chi_{2} = \frac{1}{\frac{\sin \alpha_{1} - \tan \varphi \cos \alpha_{1}}{\sin \alpha_{1} - \tan \varphi \cos \alpha_{1}}} + \frac{k_{p}}{\cos \alpha_{1} + \tan \varphi \sin \alpha_{1}}}{\chi_{1}} \\ \chi_{3} = \frac{\frac{cl_{CF} \cos \alpha_{1} + cl_{FE} \cos \alpha_{2}}{\sin \alpha_{1} - \tan \varphi \cos \alpha_{1}}}{\chi_{1}} + \frac{cl_{CF} \sin \alpha_{1} + cl_{FE} \sin \alpha_{2}}{\cos \alpha_{1} + \tan \varphi \sin \alpha_{1}}}{\chi_{1}} \\ \chi_{4} = \frac{1}{\chi_{2} - \frac{\eta_{1}k_{p}}{(\eta_{3} - \eta_{1}\eta_{2})(\cos \alpha_{1} + \tan \varphi \sin \alpha_{1})}} \end{cases}$$
(12)

于是,由式(10)可得 N<sub>pn</sub>、N<sub>n2</sub>的表达式分别为:

$$N_{pn} = \chi_{4} \begin{cases} \frac{T_{n-1}\cos\alpha_{2} - N_{n-1}\sin\alpha_{2}}{(\sin\alpha_{1} - \tan\varphi\cos\alpha_{1})\chi_{1}} + \frac{T_{n-1}\sin\alpha_{2} + N_{n-1}\cos\alpha_{2}}{(\cos\alpha_{1} + \tan\varphi\sin\alpha_{1})\chi_{1}} + \frac{G_{n}}{(\cos\alpha_{1} + \tan\varphi\sin\alpha_{1})\chi_{1}} - \frac{\eta_{1}\eta_{4}}{\eta_{3} - \eta_{1}\eta_{2}} - \chi_{3} - \frac{\eta_{1}\eta_{4}}{\eta_{3} - \eta_{1}\eta_{2}} - \chi_{3} - \frac{\eta_{1}\eta_{4}}{\eta_{3} - \eta_{1}\eta_{2}} - \chi_{3} - \frac{\eta_{1}\eta_{4}}{(\cos\alpha_{1} + \tan\varphi\sin\alpha_{1})\chi_{1}} - \frac{\eta_{1}\eta_{4}}{(\cos\alpha_{1} + \eta_{1})\chi_{1}} - \frac{\eta_{1}\eta_{1}}{(\cos\alpha_{1} + \eta_{1})\chi_{1}} - \frac{\eta_{1}\eta_{2}}{(\cos\alpha_{1} + \eta_{1})\chi_{1}} - \frac{\eta_{1}\eta_{2}}{(\cos\alpha_{1} + \eta_{1})\chi_{1}} - \frac{\eta_{1}\eta_{2}}{(\cos\alpha_{1} + \eta_{1})\chi_{1}} - \frac{\eta_{1}\eta_{2}}{(\cos\alpha_{1} + \eta_{1})\chi_{1}} - \frac{\eta_{1}\eta_{1}}{(\cos\alpha_{1} + \eta_{1})\chi_{1}} - \frac{\eta_{1}\eta_{1$$

$$N_{n2} =$$

$$\chi_{2}\chi_{4} \begin{cases} \frac{T_{n-1}\cos\alpha_{2} - N_{n-1}\sin\alpha_{2}}{(\sin\alpha_{1} - \tan\varphi\cos\alpha_{1})\chi_{1}} + \frac{T_{n-1}\sin\alpha_{2} + N_{n-1}\cos\alpha_{2}}{(\cos\alpha_{1} + \tan\varphi\sin\alpha_{1})\chi_{1}} + \frac{G_{n}}{(\cos\alpha_{1} + \tan\varphi\sin\alpha_{1})\chi_{1}} - \frac{\eta_{1}\eta_{4}}{\eta_{3} - \eta_{1}\eta_{2}} - \chi_{3} - \frac{\eta_{1}\eta_{4}}{\eta_{3} - \eta_{1}\eta_{2}} - \chi_{3} - \frac{\eta_{1}\eta_{4}}{\eta_{3} - \eta_{1}\eta_{2}} - \chi_{3} - \frac{1}{\eta_{3} - \eta_{1}\eta_{2}} \frac{T_{n-1}\sin\alpha_{2} + N_{n-1}\cos\alpha_{2} + G_{n}}{\cos\alpha_{1} + \tan\varphi\sin\alpha_{1}} - \frac{1}{\eta_{3} - \eta_{1}\eta_{2}} - \chi_{3} - \frac{\eta_{1}\eta_{4}}{\eta_{3} - \eta_{1}\eta_{2}} - \chi_{3} - \frac{\eta_{1}\eta_{4}}{(\cos\alpha_{1} + \tan\varphi\sin\alpha_{1})} - \frac{1}{\eta_{3} - \eta_{1}\eta_{2}} - \chi_{3} - \frac{\eta_{1}\eta_{4}}{\eta_{3} - \eta_{1}\eta_{2}} - \chi_{3} - \frac{\eta_{1}\eta_{4}}{(\cos\alpha_{1} + \tan\varphi\sin\alpha_{1})} + \frac{1}{\eta_{3} - \eta_{1}\eta_{2}} - \frac{1}{\eta_{3} - \eta_{1}\eta_{2}} - \frac{1}{\eta_{3} - \eta_{1}\eta_{2}} - \frac{1}{\eta_{3} - \eta_{1}\eta_{2}} - \frac{1}{(\cos\alpha_{1} + \tan\varphi\sin\alpha_{1})\chi_{1}} - \frac{1}{(\cos\alpha_{1} + \cos\alpha_{1})\chi_{1}} - \frac{1}{(\cos\alpha_{1} + \tan\varphi\sin\alpha_{1})\chi_{1}} - \frac{1}{(\cos\alpha_{1} + \cos\alpha_{1})\chi_{1}} - \frac{1}{(\cos\alpha_{1} + \cos\alpha_{1})\chi_{1}} - \frac{1}{(\cos\alpha_{1} + \cos\alpha_{1})\chi_{1}} - \frac$$

将式(13)(14)代入式(9),可确定 *T*<sub>n1</sub>、*T*<sub>n2</sub>、*N*<sub>n1</sub> 及 *T*<sub>pn</sub>。

通过上述求解方法,可得各条块对前排桩的作用 力(局部作用力),因此将其叠加可得前排桩所受的合 推力,即:

$$N_{\rm p} = \sum_{i=1}^{n} N_{\rm pi}$$
 (15)

### 2 实例分析

西南地区某铁路工程沿线存在一中风化大理岩

上覆碎石土堆积体的潜在滑坡,滑坡主轴横断面示意 图如图 4 所示,通过现场勘查与试验,确定坡体主要 物理力学参数如表 1 所示。利用传递系数法<sup>[27]</sup>求得 天然工况下该坡体稳定系数为 1.128。拟采用双排抗 滑桩加固,桩体采用 C30 混凝土浇筑,截面尺寸为 2 m×3 m,桩间距(平面外)为 5 m。前排桩拟设于距坡 脚水平距离约 30 m 的滑面平缓地段,后排桩则分别考 虑距前排桩水平距离 25 m 和 45 m 2 个位置,同时也 考虑后排桩顶沉埋情况,故拟定 4 种设桩方案。





表1 坡体及抗滑桩主要物理力学参数

Table 1	Main physical	and mechanical	parameters of the
	land	Islide and niles	

F								
材料类型	土体重度/ (kN·m <sup>-3</sup> )	黏聚力/ kPa	内摩擦角/ (°)	弹性模量/ MPa	泊松比			
碎石土	22	12	21	50	0.33			
大理岩	23	500	40	600	0.25			
抗滑桩	25	-	-	-	0.22			

方案一:后排桩位于桩位一,桩顶位于坡面;

方案二:后排桩位于桩位二,桩顶位于坡面;

方案三:后排桩位于桩位一,桩顶埋深 7.4 m;

方案四:后排桩位于桩位二,桩顶埋深7.0m。

在设计安全系数为1.20的情况下,利用传递系数 法及地基系数"*m*"法<sup>[27]</sup>求得各方案时后排桩对其前侧 滑体的推力分布如图 5 所示。

采用前述斜条分法,求解作用于前排桩受荷段后侧的坡体推力。各设桩方案相应的排间土体条分模式见图 6,前排桩土界面综合参数 k<sub>n</sub>取为 tang<sup>[28]</sup>。

表2给出了各方案时相应各土条对前排桩的局部 推力 N<sub>pi</sub> 及其合推力值。可见,对于后排全长桩的方 案一和方案二,其前排桩推力计算值比较接近,方案 一(排间直线型滑面)较方案二(排间折线型滑面)的 推力值偏大约5%。同时,当排间滑面为直线型时(方





案一和方案三),方案三(后桩埋深 7.4 m)较方案一 (后桩全长)的前排桩推力值偏小约 10.6%;当排间为



Fig. 6 Slice modes of slide mass between two piles corresponding to four design plans in the example

折线型滑面时(方案二和方案四),方案四(后桩埋深 7.0 m)较方案二(后桩全长)的前排桩推力值偏小约 8.6%。此外,所拟定的4种方案中,方案四时作用于 前排桩后侧的推力最小。方案一和方案二的前排桩 推力值均大于后排沉埋桩的方案三和方案四,这也反 映出后排桩沉埋有利于减小前排桩后侧推力。其原 因在于后排桩沉埋时作用于排间滑体后侧的推力有 所减小(图 5)。

对表 2 中所示的各土条对前排桩的局部推力值除 以其相应作用面积,可近似得到该条处前排桩后侧的 压应力,进而可得前排桩受荷段后侧推力分布模式如 图 7 所示。

由图 7 可见,各设桩方案时前排桩后侧推力分布 模式较为相似,均呈不规则的抛物线型分布。桩顶局 部范围内桩后压应力有先减小再增大的特征,最大压 应力位于距受荷段底端高 0.11~0.17 倍受荷段高度, 其下至滑面处桩后压应力急剧减小至近于零。

	表 2 任 4 种设桩万案卜各土条对刖排桩局部推刀及具合推刀	
Table 2	Local thrusts by each slice and their resultants on the fore piles under four design plans	/ $(kN \cdot m^{-1})$

				-					-			-		
况检查安	条块编号								× +6L					
₩ 反 性 力 杀 一	1	2	3	4	5	6	7	8	9	10	11	12	13	忌推力
方案一	195.10	179.20	193.60	156.20	216.80	269.00	317.20	363.20	408.30	452.90	497.30	541.50	0	3 790.40
方案二	316.60	286.60	284.40	182.20	226.00	270.90	316.50	362.50	409.00	455.60	502.40	0	-	3 612.60
方案三	213.50	196.10	198.20	146.90	198.70	241.10	279.00	314.70	349.20	383.20	416.90	450.40	0	3 387.90
方案四	334.20	289.10	271.30	166.10	202.60	240.50	279.40	319.00	359.10	399.40	440.00	0	_	3 300.70



Fig. 7 Distribution curves of thrust on the fore piles under four design plans in the example

## 3 对比验证

为了进一步说明本文方法的合理性,以前述实例中的设桩方案四为例,采用FLAC<sup>3D</sup>进行数值模拟分析,与理论计算结果予以比较。所建数值模型如图 8 所示,模型含 72 198 个 8 节点 6 面体单元,坡体采用服从 Mohr-Coulomb 屈服准则和关联流动法则的理想 弹塑性本构模型,桩体视为弹性材料,利用结构单元 模拟。模型前后左右四个边界进行水平位移约束,底面同时进行水平和竖向位移约束。采用强度折减法 进行数值模拟分析,得到该滑坡模型的稳定系数为 1.12,与前述的传递系数法计算结果基本一致,说明数 值模型有一定的合理性。





图 9 为数值模拟得到的加固后滑坡的水平应力分 布云图,由此可得前排桩后压应力分布图,如图 10 所 示。可见,在桩顶以下局部约 0.25 倍受荷段高度范围 以内,理论计算与FLAC<sup>3D</sup>模拟所得的桩后压应力分 布有差异,前者沿深度呈非线性减小模式,后者则为 近似线性增大模式,且后者整体小于前者,意味着理 论计算偏于保守;在此局部范围以下,两种方法得到 的桩后坡体压力变化规律基本一致,均为先近似线性 增大,近滑面处达到极大值,然后再急剧减小;二者得 到的桩后峰值压应力偏差约10%,理论算法略小于数 值模拟结果。









表3给出了传递系数法、FLAC<sup>3D</sup>模拟法和本文的斜条分法得到的4种设桩方案时的前排桩受荷段后侧坡体推力值。可见,本文方法比传递系数法的结果分别偏大约10.6%、12.5%、11.5%、13.4%,而比FLAC<sup>3D</sup>结果分别偏大约13.2%、7.95%、1.31%、-0.72%。因此,相对于传递系数法,本文方法整体更接近于FLAC<sup>3D</sup>计算结果。

表 3	不同方法求得的 4 种i	殳桩方案↑	「前排胡	E上推力值对比			
	Table 3 Comp	parison of	thrust o	)n			
the fore piles under four design							

	plans by	y various me	tnods	(KIN'III )
计算方法	方案一	方案二	方案三	方案四
斜条分法	3 790.36	3 612.60	3 387.90	3 300.70
传递系数法	3 425.60	3 212.00	3 039.70	2 910.00
FLAC <sup>3D</sup>	3 349.03	3 346.40	3 344.10	3 324.50

### 4 结论

(1)前排桩后侧推力呈不规则的抛物线形分布, 其峰值点位置接近滑面,在峰值点以下,受荷段桩后 坡体压力急剧减小至接近于零。

(2)对于后排桩沉埋模式的双排桩,前排桩后侧 坡体推力小于后排桩全长模式的结果。本文方法实 例计算表明,前者比后者偏小8.6%~10.6%。

(3)本文提出的斜条分法与传递系数法、FLAC<sup>3D</sup>数值模拟方法计算的前排桩后侧坡体推力值较为相近,本文方法整体较既有两种方法均偏大,结果偏于保守一面。实例分析表明,本文方法比传递系数法结果偏大 10.6%~13.4%,且相对更接近于 FLAC<sup>3D</sup>数值模拟结果。

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